EG 2020 STAR presentation



Survey of Models for Acquiring the Optical Properties of Translucent Materials

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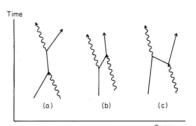
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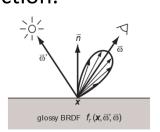




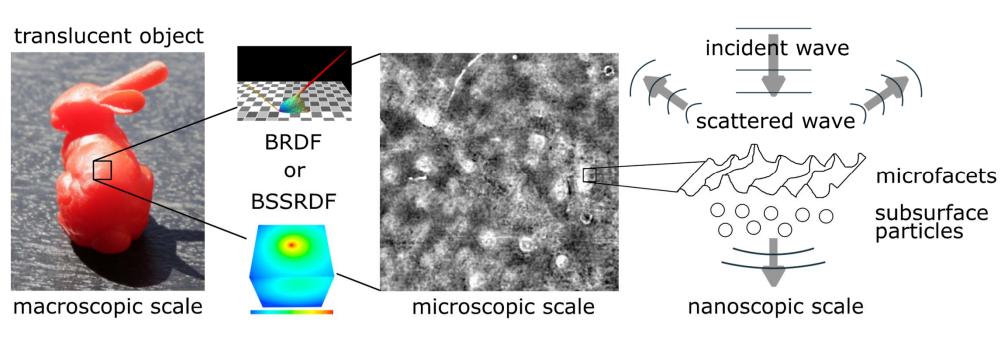
Optical properties

- Parameters that determine how light interacts with a material.
- Quantum and wave theories:
 - Quantum scale: photon-electron interactions in atomic systems.
 - Nanoscopic scale: charge and current densities in atomic systems.
 - Microscopic scale: polarisation and magnetisation vectors.
 - Macroscopic scale: permittivity, permeability, conductivity.
- Radiative transfer theory:
 - Microscopic scale: complex index of refraction.
 - Mesoscopic scale: surface BSDF, scattering cross section, phase function.
 - Macroscopic scale: scattering properties, BSSRDF, BRDF, BTDF.





Multiscale modelling



• With simulation of light propagation, we can compute macroscopic optical properties by considering geometry at different scales.

Index of refraction (or refractive index)

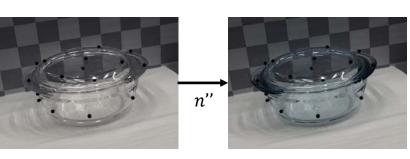
• Combining permittivity (ϵ), permeability (μ), and conductivity (σ):

•
$$n_{\text{med}} = n' + i n'' = c \sqrt{\mu \left(\varepsilon + i \frac{\sigma}{\omega}\right)}$$

- ω is angular frequency.
- *c* is the speed of light *in vacuo*.
- Real part $n' \approx \frac{c}{v}$
 - v is the phase velocity of the light wave.
- Imaginary part $n'' \approx \frac{\sigma_a \lambda}{4\pi}$
 - σ_a is the absorption coefficient.
 - λ is the wavelength *in vacuo*.



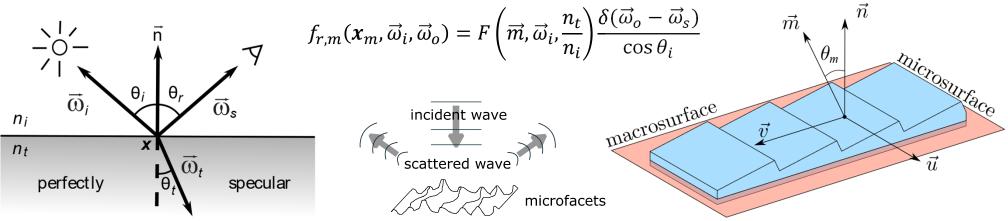
varying the real part n'



Including absorption

Microfacet BSDF

- A surface is **optically smooth** if the surface roughness R_q is sufficiently small compared with the wavelength λ .
- Rayleigh smooth-surface criterion: $R_a < \lambda/(8\cos\theta_i)$.
- ullet Considering smooth microgeometry we can use $n_{
 m med}$ as input for analytic or computational solutions for Maxwell's equations.
- Example: Fresnel reflectance F for a microfacet BSDF.



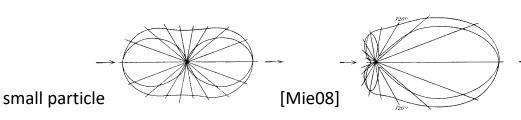
Particle phase function and cross sections

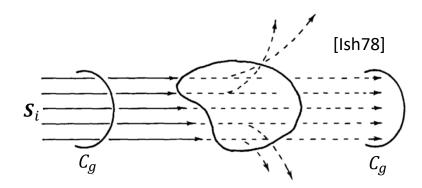
Particle cross sections

- C_g is the geometric cross section.
- C_s is the scattering cross section.
- C_a is the absorption cross section.
- $C_t = C_s + C_a$ is the extinction cross section.



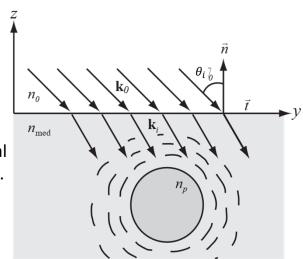
- $p_m(\vec{\omega}_i, \vec{\omega}_o)$ is the far field distribution of the scattered light.
- $g = \int_{4\pi} p_m(\vec{\omega}_i, \vec{\omega}_o) (\vec{\omega}_i \cdot \vec{\omega}_o) d\omega$ is the asymmetry parameter in [-1,1].





Example: Insert $x=\frac{2\pi r n_{\mathrm{med}}}{\lambda}$ and $y=\frac{2\pi r n_p}{\lambda}$ in Lorenz-Mie theory to compute C_s , C_t , and p of a spherical particle of radius r.

large particle



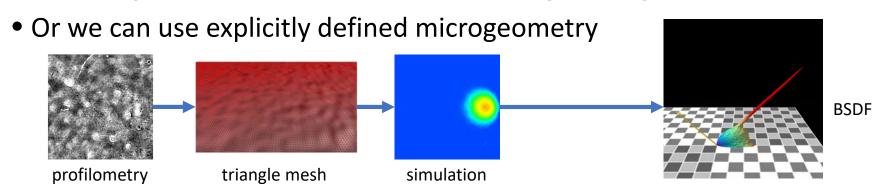
Scattering properties of a medium

[FCJ07]

- Using a particle size distribution N(r): $\sigma_s = \int_{r_{\min}}^{r_{\max}} C_s(r) N(r) \, \mathrm{d}r$ • σ_s is the scattering coefficient.
 - Similarly for σ_a (absorption coefficient) and p (ensemble phase function).
- Using a microfacet normal distribution $D(\vec{m})$:

$$f_{S}(\vec{\omega}_{i}, \vec{\omega}_{o}, \vec{n}) = \int \left| \frac{\vec{\omega}_{i} \cdot \vec{m}}{\vec{\omega}_{i} \cdot \vec{n}} \right| f_{m}(\vec{\omega}_{i}, \vec{\omega}_{o}, \vec{m}) \left| \frac{\vec{\omega}_{o} \cdot \vec{m}}{\vec{\omega}_{o} \cdot \vec{n}} \right| G(\vec{\omega}_{i}, \vec{\omega}_{o}, \vec{m}) D(\vec{m}) d\omega_{m}$$

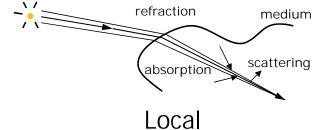
• G is a geometrical attenuation term (shadowing/masking).

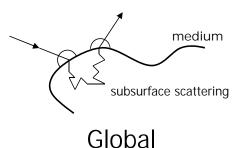


Global scattering function (BSSRDF)

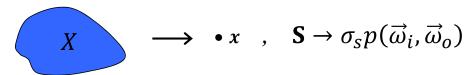
• From local to global formulation using scattering operators [Pre65].

$$L_r = L_i \mathbf{S} = L_i \mathbf{F}_s \sum_{j=0}^{\infty} \mathbf{S}^j \mathbf{F}_s$$



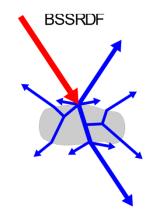


- **F**_s is surface scattering
- **S**⁰ is direct transmission
- S^{j} for j > 0 is subsurface scattering (with j scattering events)
- For



• Continuous boundary and interior leads to the BSSRDF:

$$S(X; \boldsymbol{x}_i, \vec{\omega}_i; \boldsymbol{x}_o, \vec{\omega}_o) = \lim_{\substack{X_i \to \boldsymbol{x}_i \\ \Omega_i \to \vec{\omega}_i}} \frac{L_i \mathbf{S}(\boldsymbol{x}_o, \vec{\omega}_o)}{L_i(X_i, \Omega_i) A_{i\perp}(X_i) \omega_i(\Omega_i)} = \frac{\mathrm{d}L_r(\boldsymbol{x}_o, \vec{\omega}_o)}{L_i(\boldsymbol{x}_i, \vec{\omega}_i) \mathrm{d}A_{i\perp} \, \mathrm{d}\omega_i} = \frac{\mathrm{d}L_r(\boldsymbol{x}_o, \vec{\omega}_o)}{\mathrm{d}\Phi_i(\boldsymbol{x}_i, \vec{\omega}_i)}$$



Macroscopic BRDF/BTDF

- Object with homogeneous scattering properties.
- Uniform irradiation of the object over an area A_i around x_o .
 - A_i is large enough to include all x_i with subsurface scattering to x_o .
- The BRDF is then

$$f_r(\mathbf{x}, \vec{\omega}_i, \vec{\omega}_o) = \int_{A_i} S(X; \mathbf{x}_i, \vec{\omega}_i; \mathbf{x}_o, \vec{\omega}_o) \, dA_i = \frac{dL_r(\mathbf{x}, \vec{\omega}_o)}{dE(\mathbf{x}, \vec{\omega}_i)}$$

- The equation is the same for the BTDF, but then $\vec{\omega}_i \cdot \vec{n}_o < 0$.
 - \vec{n}_o is the surface normal at the point of observation.
- Macroscopic BRDF/BTDF works well for opaque/thin objects.
- Not a good approximation for solid translucent objects.



Lambertian BRDF approximation

Appearance of translucent materials

• Varying the transport mean free path $\frac{1}{\sigma_t'} = \frac{1}{\sigma_a + (1-g)\sigma_s}$



- Varying surface roughness
- Varying lighting environment









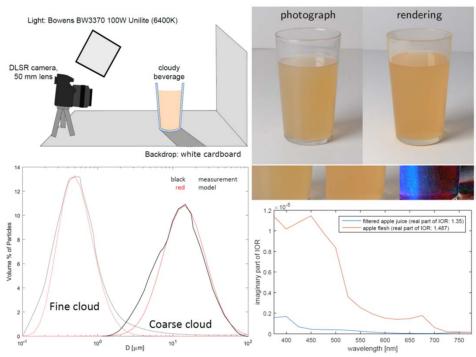


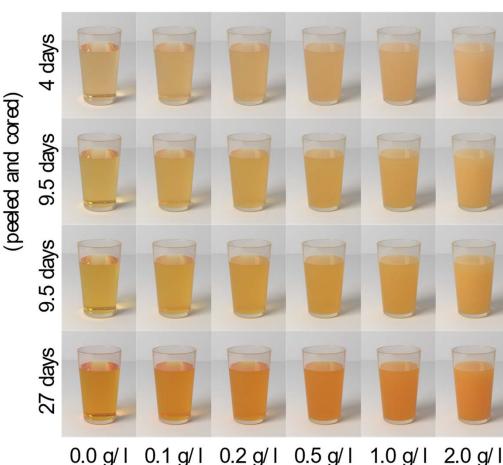
Influence of particle content

[DFKB16]

- Apple juice example
 - Particle concentration (horizontally).
 - Storage time and handling (vertically).

[FK19]





Discussion

- Translucent objects require optical properties describing both surface and subsurface scattering.
- What is the best appearance specification for translucent objects?

Influence of surface roughness [LFD*20]

